# CERTAIN CLASSES OF BI-UNIVALENT FUNCTIONS ASSOCIATED WITH $q$-ANALOGUE OF BESSEL FUNCTIONS 

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Abstract: In this paper we consider various subclasses of bi-univalent functions defined by the Horadam polynomials associated with $q$-analogue of Bessel functions. Further, we obtain coefficient estimates and Fekete-Szegö inequalities for the defined classes.

Keywords and Phrases: Univalent functions, bi-univalent functions, bi-convex functions, bi-starlike functions, Fekete-Szegö inequality, $q$-derivative operator, Horadam polynomials, Bessel functions.

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## 1. Introduction

Let $\mathcal{A}$ denote the class of functions of the form

$$
\begin{equation*}
f(z)=z+\sum_{n=2}^{\infty} a_{n} z^{n} \tag{1.1}
\end{equation*}
$$

which are analytic in the open unit disk $\Delta=\{z: z \in \mathbb{C}$ and $|z|<1\}$. Further, by $\mathcal{S}$ we shall denote the class of all functions in $\mathcal{A}$ which are univalent in $\Delta$.

It is well known that every function $f \in \mathcal{S}$ has an inverse $f^{-1}$, defined by

$$
f^{-1}(f(z))=z \quad(z \in \Delta)
$$

and

$$
f\left(f^{-1}(w)\right)=w \quad\left(|w|<r_{0}(f) ; r_{0}(f) \geqq \frac{1}{4}\right),
$$

where

$$
\begin{equation*}
f^{-1}(w)=w-a_{2} w^{2}+\left(2 a_{2}^{2}-a_{3}\right) w^{3}-\left(5 a_{2}^{3}-5 a_{2} a_{3}+a_{4}\right) w^{4}+\cdots . \tag{1.2}
\end{equation*}
$$

A function $f \in \mathcal{A}$ is said to be bi-univalent in $\Delta$ if both the function $f$ and its inverse $f^{-1}$ are univalent in $\Delta$. Let $\sigma$ denote the class of bi-univalent functions in $\Delta$ given by (1.1).

In 2010, Srivastava et al. [35] revived the study of bi-univalent functions by their pioneering work on the study of coefficient problems. Various subclasses of the bi-univalent function class $\sigma$ were introduced and non-sharp estimates on the first two coefficients $\left|a_{2}\right|$ and $\left|a_{3}\right|$ in the Taylor-Maclaurin series expansion (1.1) were found in the very recent investigations (see, for example, [1], [2], [3], [4], [5], [6], [7], [8], [9], [10], [13], [14], [16], [19], [20], [22], [23], [25], [26], [27], [28], [29], [31], [32], [33], [34], [36], [37]) and including the references therein. The afore-cited all these papers on the subject were actually motivated by the work of Srivastava et al. [35]. However, the problem to find the coefficient bounds on $\left|a_{n}\right|(n=3,4, \cdots)$ for functions $f \in \sigma$ is still an open problem.

For analytic functions $f$ and $g$ in $\Delta, f$ is said to be subordinate to $g$ if there exists an analytic function $w$ such that (see, for example, [12], [24])

$$
w(0)=0, \quad|w(z)|<1 \quad \text { and } \quad f(z)=g(w(z)) .
$$

This subordination will be denoted here by

$$
f \prec g
$$

or, conventionally, by

$$
f(z) \prec g(z) .
$$

In particular, when $g$ is univalent in $\Delta$,

$$
f \prec g \quad(z \in \Delta) \Leftrightarrow f(0)=g(0) \quad \text { and } \quad f(\Delta) \subset g(\Delta) .
$$

The Bessel function of the first kind of order $\nu$ is defined by the infinite series (see [21])

$$
\begin{equation*}
J_{\nu}(z):=\sum_{n=0}^{\infty} \frac{(-1)^{n}(z / 2)^{2 n+\nu}}{n!\Gamma(n+\nu+1)}, \quad(z \in \mathbb{C}, \nu \in \mathbb{R}) \tag{1.3}
\end{equation*}
$$

where $\Gamma$ stands for the Gamma function. Recently, Szasz and Kupan [30] investigated the univalence of the normalized Bessel function of the first kind $\kappa_{\nu}: \Delta \rightarrow \mathbb{C}$ defined by

$$
\begin{align*}
\kappa_{\nu}(z) & :=2^{\nu} \Gamma(\nu+1) z^{1-\nu / 2} J_{\nu}\left(z^{1 / 2}\right) \\
& =z+\sum_{n=2}^{\infty} \frac{(-1)^{n-1} \Gamma(\nu+1)}{4^{n-1}(n-1)!\Gamma(n+\nu)} z^{n}, \quad(z \in \Delta, \nu \in \mathbb{R}) \tag{1.4}
\end{align*}
$$

For $0<q<1$, El-Deeb and Bulboaca [15] defined the $q$-derivateive operator for $\kappa_{\nu}$ as follows:

$$
\begin{align*}
\partial_{q} \kappa_{\nu}(z) & =\partial_{q}\left[z+\sum_{n=2}^{\infty} \frac{(-1)^{n-1} \Gamma(\nu+1)}{4^{n-1}(n-1)!\Gamma(n+\nu)} z^{n}\right]:=\frac{\kappa_{\nu}(q z)-\kappa_{\nu}(z)}{z(q-1)} \\
& =1+\sum_{n=2}^{\infty} \frac{(-1)^{n-1} \Gamma(\nu+1)}{4^{n-1}(n-1)!\Gamma(n+\nu)}[n, q] z^{n-1}, \quad(z \in \Delta) \tag{1.5}
\end{align*}
$$

where

$$
\begin{equation*}
[n, q]:=\frac{1-q^{n}}{1-q}=1+\sum_{j=1}^{n-1} q^{j}, \quad[0, q]:=0 \tag{1.6}
\end{equation*}
$$

Using (1.6), we will define the next two products:

1. For any nonnegative integer $n$, the $q$-shifted factorial is given by

$$
[n, q]:= \begin{cases}1, & \text { if } n=0  \tag{1.7}\\ {[1, \mathrm{q}][2, \mathrm{q}] \ldots[\mathrm{k}, \mathrm{q}]} & \text { if } n \in \mathbb{N} .\end{cases}
$$

2. For any positive number $r$, the $q$ - generalized Pochhammer symbol is defined by

$$
[r, q]_{n}:= \begin{cases}1, & \text { if } n=0  \tag{1.8}\\ {[\mathrm{r}, \mathrm{q}][\mathrm{r}+1, \mathrm{q}] \ldots[\mathrm{r}+\mathrm{k}-1, \mathrm{q}]} & \text { if } n \in \mathbb{N} .\end{cases}
$$

For $\nu>0, \lambda>-1$ and $0<q<1$, El-Deeb and Bulboacǎ [15] defined the function $\mathcal{J}_{\nu, q}^{\lambda}: \Delta \rightarrow \mathbb{C}$ by (see [14], [16])

$$
\begin{equation*}
\mathcal{J}_{\nu, q}^{\lambda}(z):=z+\sum_{n=2}^{\infty} \frac{(-1)^{n-1} \Gamma(\nu+1)}{4^{n-1}(n-1)!\Gamma(n+\nu)} \frac{[n, q]!}{[\lambda+1, q]_{n-1}} z^{n}, \quad(z \in \Delta) \tag{1.9}
\end{equation*}
$$

A simple computation shows that

$$
\begin{equation*}
\mathcal{J}_{\nu, q}^{\lambda}(z) * \mathcal{M}_{q, \lambda+1}(z)=z \partial_{q} \kappa_{\nu}(z), \quad(z \in \Delta) \tag{1.10}
\end{equation*}
$$

where the function $\mathcal{M}_{q, \lambda+1}(z)$ is given by

$$
\begin{equation*}
\mathcal{M}_{q, \lambda+1}(z):=z+\sum_{n=2}^{\infty} \frac{[\lambda+1, q]_{n-1}}{[n-1, q]!} z^{n}, \quad(z \in \Delta) \tag{1.11}
\end{equation*}
$$

Using the definition of $q$-derivative along with the idea of convolutions, El-Deeb and Bulboacǎ [15] introduced the linear operator $\mathcal{N}_{\nu, q}^{\lambda}: \mathcal{A} \rightarrow \mathcal{A}$ defined by

$$
\begin{align*}
\mathcal{N}_{\nu, q}^{\lambda} f(z) & :=\mathcal{J}_{\nu, q}^{\lambda} * f(z) \\
& =z+\sum_{n=2}^{\infty} \psi_{n} a_{n} z^{n}, \quad(\nu>0, \lambda>-1,0<q<1, z \in \Delta) \tag{1.12}
\end{align*}
$$

where

$$
\begin{equation*}
\psi_{n}:=\frac{(-1)^{n-1} \Gamma(\nu+1)}{4^{n-1}(n-1)!\Gamma(n+\nu)} \frac{[n, q]!}{[\lambda+1, q]_{n-1}} \tag{1.13}
\end{equation*}
$$

Remark 1.1. [15] From the definition relation (1.12), we can easily verify that the next relations hold for all $f \in \mathcal{A}$ :

$$
\begin{equation*}
[\lambda+1, q] \mathcal{N}_{\nu, q}^{\lambda} f(z)=[\lambda, q] \mathcal{N}_{\nu, q}^{\lambda+1} f(z)+q^{\lambda} z \partial_{q}\left([\lambda+1, q] \mathcal{N}_{\nu, q}^{\lambda+1} f(z)\right), \quad z \in \Delta \tag{1.14}
\end{equation*}
$$

and

$$
\begin{align*}
\lim _{q \rightarrow 1^{-}} \mathcal{N}_{\nu, q}^{\lambda} f(z) & =\mathcal{J}_{\nu, 1}^{\lambda} f(z):=\mathcal{J}_{\nu}^{\lambda} f(z) \\
& =z+\sum_{n=2}^{\infty} \frac{(-1)^{n-1} \Gamma(\nu+1)}{4^{n-1}(n-1)!\Gamma(n+\nu)} \frac{n!}{(\lambda+1)_{n-1}} a_{n} z^{n}, \quad(z \in \Delta) \tag{1.15}
\end{align*}
$$

The Horadam polynomials $h_{n}(x, a, b ; p, q)$, or briefly $h_{n}(x)$ are given by the following recurrence relation (see [17], [18])):

$$
\begin{equation*}
h_{1}(x)=a, \quad h_{2}(x)=b x \quad \text { and } \quad h_{n}(x)=\quad p x h_{n-1}(x)+q h_{n-2}(x) \quad(n \geq 3) \tag{1.16}
\end{equation*}
$$

for some real constants $a, b, p$ and $q$.
The generating function of the Horadam polynomials $h_{n}(x)$ (see [18]) is given by

$$
\begin{equation*}
\Pi(x, z):=\sum_{n=1}^{\infty} h_{n}(x) z^{n-1}=\frac{a+(b-a p) x z}{1-p x z-q z^{2}}, \quad 1-p x z-q z^{2} \neq 0, \quad \forall z \in \Delta . \tag{1.17}
\end{equation*}
$$

Here, and in what follows, the argument $x \in \mathbb{R}$ is independent of the argument $z \in \mathbb{C}$; that is, $x \neq \Re(z)$.

Note that for particular values of $a, b, p$ and $q$, the Horadam polynomial $h_{n}(x)$ leads to various polynomials, among those, we list a few cases here (see, [17], [18] for more details):

1. For $a=b=p=q=1$, we have the Fibonacci polynomials $F_{n}(x)$.
2. For $a=2$ and $b=p=q=1$, we obtain the Lucas polynomials $L_{n}(x)$.
3. For $a=q=1$ and $b=p=2$, we get the Pell polynomials $P_{n}(x)$.
4. For $a=b=p=2$ and $q=1$, we attain the Pell-Lucas polynomials $Q_{n}(x)$.
5. For $a=b=1, p=2$ and $q=-1$, we have the Chebyshev polynomials $T_{n}(x)$ of the first kind
6. For $a=1, b=p=2$ and $q=-1$, we obtain the Chebyshev polynomials $U_{n}(x)$ of the second kind.

Abirami et al. [1] considered bi- Mocanu - convex functions and bi- $\mu$ - starlike functions to discuss initial estimations of Taylor-Macularin series which is associated with Horadam polynomials, Abirami et al. [2] discussed coefficient estimates for the classes of $\lambda$-bi-pseudo-starlike and bi-Bazilevič functions using Horadam polynomial, Alamoush [3], [4] defined subclasses of bi-starlike and bi-convex functions involving the Poisson distribution series involving Horadam polynomials and a class of bi-univalent functions associated with Horadam polynomials respectively and obtained initial coefficient estimates, Altınkaya and Yalçın [7], [8] obtained coefficient estimates for Pascu-type bi-univalent functions and for the class of linear combinations of bi-univalent functions by means of $(p, q)$-Lucas polynomials respectively, Aouf et al. [10] discussed initial coefficient estimates for general class of pascu-type bi-univalent functions of complex order defined by $q$-Sălăgean operator and associated with Chebyshev polynomials, Awolere and Oladipo [11] found initial
coefficients of bi-univalent functions defined by sigmoid functions involving pseudostarlikeness associated with Chebyshev polynomials, Naeem et al. [22] considered a general class of bi-Bazilevič type functions associated with Faber polynomial to discuss $n$-th coefficients estimates, Magesh and Bulut [23] discussed Chebyshev polynomial coefficient estimates for a class of analytic bi-univalent functions related to pseudo-starlike functions, Orhan et al. [25] discussed initial estimates and Fekete-Szegö bounds for bi-Bazilevič functions related to shell-like curves, Sakar and Aydogan [28] obtained initial bounds for the class of generalized Sălăgean type bi- $\alpha-$ convex functions of complex order associated with the Horadam polynomials, Singh et al. [31] found coefficient estimates for bi- $\alpha$-convex functions defined by generalized Sãlãgean operator related to shell-like curves connected with Fibonacci numbers, Srivastava et al. [32] introduced a technique by defining a new class of bi-univalent functions associated with the Horadam polynomials to discuss the coefficient estimates, Srivastava et al. [34] gave a direction to study the Faber polynomial coefficient estimates for bi-univalent functions defined by the Tremblay fractional derivative operator, Srivastava et al. [36] obtained general coefficient $\left|a_{n}\right|$ for a general class analytic and bi-univalent functions defined by using differential subordination and a certain fractional derivative operator associated with Faber polynomial, Wanas and Alina [37] discussed applications of Horadam polynomials on Bazilevič bi-univalent functions by means of subordination and found initial bounds. Motivated in these lines, estimates on initial coefficients of the TaylorMaclaurin series expansion (1.1) and Fekete-Szegö inequalities for certain classes of bi-univalent functions defined by means of Horadam polynomials are obtained. The classes introduced in this paper are motivated by the corresponding classes investigated in [19], [14].

## 2. Coefficient Estimates and Fekete-Szegö Inequalities

Definition 2.1. A function $f \in \mathcal{A}$ of the form (1.1) belongs to the class $\mathcal{M}_{\sigma}(\alpha$, $\lambda, \nu, q, x)$ for $\alpha \geq 0, \nu>0, \lambda>-1,0<q<1$, and $z, w \in \Delta$, if the following conditions are satisfied:

$$
\alpha\left(1+\frac{z\left(\mathcal{N}_{\nu, q}^{\lambda} f(z)\right)^{\prime \prime}}{\left(\mathcal{N}_{\nu, q}^{\lambda} f(z)\right)^{\prime}}\right)+(1-\alpha) \frac{1}{\left(\mathcal{N}_{\nu, q}^{\lambda} f(z)\right)^{\prime}} \prec \Pi(x, z)+1-a
$$

and for $g$ the analytic extension (continuation) of $f^{-1}$ given by (1.2)

$$
\alpha\left(1+\frac{w\left(\mathcal{N}_{\nu, q}^{\lambda} g(w)\right)^{\prime \prime}}{\left(\mathcal{N}_{\nu, q}^{\lambda} g(w)\right)^{\prime}}\right)+(1-\alpha) \frac{1}{\left(\mathcal{N}_{\nu, q}^{\lambda} g(w)\right)^{\prime}} \prec \Pi(x, w)+1-a
$$

where the real constant $a$ is as in (1.16).

Remark 2.1. Putting $q \rightarrow 1^{-}$, we obtain that

$$
\lim _{q \rightarrow 1^{-}} \mathcal{M}_{\sigma}(\alpha, \lambda, \nu, q, x)=: \mathcal{M}_{\sigma}(\alpha, \lambda, \nu, x),
$$

where $f \in \sigma$,

$$
\alpha\left(1+\frac{z\left(\mathcal{J}_{\nu}^{\lambda} f(z)\right)^{\prime \prime}}{\left(\mathcal{J}_{\nu}^{\lambda} f(z)\right)^{\prime}}\right)+(1-\alpha) \frac{1}{\left(\mathcal{J}_{\nu}^{\lambda} f(z)\right)^{\prime}} \prec \Pi(x, z)+1-a
$$

and for $g$ the analytic extension (continuation) of $f^{-1}$ given by (1.2)

$$
\alpha\left(1+\frac{w\left(\mathcal{J}_{\nu}^{\lambda} g(w)\right)^{\prime \prime}}{\left(\mathcal{J}_{\nu}^{\lambda} g(w)\right)^{\prime}}\right)+(1-\alpha) \frac{1}{\left(\mathcal{J}_{\nu}^{\lambda} g(w)\right)^{\prime}} \prec \Pi(x, w)+1-a,
$$

where the real constant $a$ is as in (1.16).
For functions in the class $\mathcal{M}_{\sigma}(\alpha, \lambda, \nu, q, x)$, the following coefficient estimates and Fekete-Szegö inequality are obtained.
Theorem 2.1. Let $f(z)=z+\sum_{n=2}^{\infty} a_{n} z^{n}$ be in the class $\mathcal{M}_{\sigma}(\alpha, \lambda, \nu, q, x)$. Then

$$
\begin{aligned}
& \left|a_{2}\right| \leq \frac{|b x| \sqrt{|b x|}}{\sqrt{\left|(9 \alpha-3) \psi_{3} b^{2} x^{2}-\left[(8 \alpha-4) b^{2} x^{2}+4\left(p x^{2} b+q a\right)(2 \alpha-1)^{2}\right] \psi_{2}^{2}\right|}}, \\
& \left|a_{3}\right| \leq \frac{|b x|}{(9 \alpha-3) \psi_{3}}+\frac{b^{2} x^{2}}{4(2 \alpha-1)^{2} \psi_{2}^{2}}
\end{aligned}
$$

and for $\mu \in \mathbb{R}$

$$
\begin{aligned}
& \left|a_{3}-\mu a_{2}^{2}\right| \\
& \leq\left\{\begin{array}{l}
\frac{|b x|}{(9 \alpha-3) \psi_{3}} \\
\text { if }|\mu-1| \leq \frac{\left|(9 \alpha-3) \psi_{3} b^{2} x^{2}-\left[(8 \alpha-4) b^{2} x^{2}+4\left(p x^{2} b+q a\right)(2 \alpha-1)^{2}\right] \psi_{2}^{2}\right|}{b^{2} x^{2}(9 \alpha-3) \psi_{3}} \\
\frac{|b x|^{3}|\mu-1|}{\left|(9 \alpha-3) \psi_{3} b^{2} x^{2}-\left[(8 \alpha-4) b^{2} x^{2}+4\left(p x^{2} b+q a\right)(2 \alpha-1)^{2}\right] \psi_{2}^{2}\right|} \\
\text { if }|\mu-1| \geq \frac{\left|(9 \alpha-3) \psi_{3} b^{2} x^{2}-\left[(8 \alpha-4) b^{2} x^{2}+4\left(p x^{2} b+q a\right)(2 \alpha-1)^{2}\right] \psi_{2}^{2}\right|}{b^{2} x^{2}(9 \alpha-3) \psi_{3}} .
\end{array}\right.
\end{aligned}
$$

Proof. Let $f \in \mathcal{M}_{\sigma}(\alpha, \lambda, \nu, q, x)$ be given by the Taylor-Maclaurin expansion (1.1). Then, there are analytic functions $r(z)$ and $s(w)$ such that

$$
r(0)=0 ; \quad s(0)=0, \quad\left|r_{n}\right|<1 \quad \text { and } \quad\left|s_{n}\right|<1 \quad(\forall z, w \in \Delta)
$$

and we can write

$$
\begin{equation*}
\alpha\left(1+\frac{z\left(\mathcal{N}_{\nu, q}^{\lambda} f(z)\right)^{\prime \prime}}{\left(\mathcal{N}_{\nu, q}^{\lambda} f(z)\right)^{\prime}}\right)+(1-\alpha) \frac{1}{\left(\mathcal{N}_{\nu, q}^{\lambda} f(z)\right)^{\prime}}=\Pi(x, r(z))+1-a \tag{2.1}
\end{equation*}
$$

and

$$
\begin{equation*}
\alpha\left(1+\frac{w\left(\mathcal{N}_{\nu, q}^{\lambda} g(w)\right)^{\prime \prime}}{\left(\mathcal{N}_{\nu, q}^{\lambda} g(w)\right)^{\prime}}\right)+(1-\alpha) \frac{1}{\left(\mathcal{N}_{\nu, q}^{\lambda} g(w)\right)^{\prime}}=\Pi(x, s(w))+1-a \tag{2.2}
\end{equation*}
$$

Equivalently,

$$
\begin{align*}
\alpha\left(1+\frac{z\left(\mathcal{N}_{\nu, q}^{\lambda} f(z)\right)^{\prime \prime}}{\left(\mathcal{N}_{\nu, q}^{\lambda} f(z)\right)^{\prime}}\right)+(1-\alpha) \frac{1}{\left(\mathcal{N}_{\nu, q}^{\lambda} f(z)\right)^{\prime}} \\
=1+h_{1}(x)-a+h_{2}(x) r(z)+h_{3}(x)[r(z)]^{2}+\cdots \tag{2.3}
\end{align*}
$$

and

$$
\begin{align*}
& \alpha\left(1+\frac{w\left(\mathcal{N}_{\nu, q}^{\lambda} g(w)\right)^{\prime \prime}}{\left(\mathcal{N}_{\nu, q}^{\lambda} g(w)\right)^{\prime}}\right)+(1-\alpha) \frac{1}{\left(\mathcal{N}_{\nu, q}^{\lambda} g(w)\right)^{\prime}} \\
& =1+h_{1}(x)-a+h_{2}(x) s(w)+h_{3}(x)[s(w)]^{2}+\cdots \tag{2.4}
\end{align*}
$$

From (2.3) and (2.4) and in view of (1.17), we obtain

$$
\begin{align*}
\alpha\left(1+\frac{z\left(\mathcal{N}_{\nu, q}^{\lambda} f(z)\right)^{\prime \prime}}{\left(\mathcal{N}_{\nu, q}^{\lambda} f(z)\right)^{\prime}}\right)+(1-\alpha) \frac{1}{\left(\mathcal{N}_{\nu, q}^{\lambda} f(z)\right)^{\prime}} \\
=1+h_{2}(x) r_{1} z+\left[h_{2}(x) r_{2}+h_{3}(x) r_{1}^{2}\right] z^{2}+\cdots \tag{2.5}
\end{align*}
$$

and

$$
\begin{align*}
\alpha\left(1+\frac{w\left(\mathcal{N}_{\nu, q}^{\lambda} g(w)\right)^{\prime \prime}}{\left(\mathcal{N}_{\nu, q}^{\lambda} g(w)\right)^{\prime}}\right)+(1-\alpha) \frac{1}{\left(\mathcal{N}_{\nu, q}^{\lambda} g(w)\right)^{\prime}} \\
=1+h_{2}(x) s_{1} w+\left[h_{2}(x) s_{2}+h_{3}(x) s_{1}^{2}\right] w^{2}+\cdots \tag{2.6}
\end{align*}
$$

If

$$
r(z)=\sum_{n=1}^{\infty} r_{n} z^{n} \quad \text { and } \quad s(w)=\sum_{n=1}^{\infty} s_{n} w^{n},
$$

then it is well known that

$$
\left|r_{n}\right| \leq 1 \quad \text { and } \quad\left|s_{n}\right| \leq 1 \quad(n \in \mathbb{N}) .
$$

Thus upon comparing the corresponding coefficients in (2.5) and (2.6), we have

$$
\begin{align*}
2 \psi_{2}(2 \alpha-1) a_{2} & =h_{2}(x) r_{1}  \tag{2.7}\\
(9 \alpha-3) \psi_{3} a_{3}-4(2 \alpha-1) \psi_{2}^{2} a_{2}^{2} & =h_{2}(x) r_{2}+h_{3}(x) r_{1}^{2}  \tag{2.8}\\
-2(2 \alpha-1) \psi_{2} a_{2} & =h_{2}(x) s_{1} \tag{2.9}
\end{align*}
$$

and

$$
\begin{equation*}
\left[(18 \alpha-6) \psi_{3}-4(2 \alpha-1) \psi_{2}^{2}\right] a_{2}^{2}-3(3 \alpha-1) \psi_{3} a_{3}=h_{2}(x) s_{2}+h_{3}(x) s_{1}^{2} . \tag{2.10}
\end{equation*}
$$

From (2.7) and (2.9), we can easily see that

$$
\begin{equation*}
r_{1}=-s_{1}, \quad \text { provided } \quad h_{2}(x)=b x \neq 0 \tag{2.11}
\end{equation*}
$$

and

$$
\begin{align*}
8 a_{2}^{2}(2 \alpha-1)^{2} \psi_{2}^{2} & =\left(h_{2}(x)\right)^{2}\left(r_{1}{ }^{2}+s_{1}{ }^{2}\right) \\
a_{2}^{2} & =\frac{\left(h_{2}(x)\right)^{2}\left(r_{1}{ }^{2}+s_{1}{ }^{2}\right)}{8(2 \alpha-1)^{2} \psi_{2}^{2}} . \tag{2.12}
\end{align*}
$$

If we add (2.8) to (2.10), we get

$$
\begin{equation*}
\left((18 \alpha-6) \psi_{3}-2(8 \alpha-4) \psi_{2}^{2}\right) a_{2}^{2}=\left(r_{2}+s_{2}\right) h_{2}(x)+h_{3}(x)\left(r_{1}^{2}+s_{1}^{2}\right) . \tag{2.13}
\end{equation*}
$$

By substituting (2.12) in (2.13), we obtain

$$
\begin{equation*}
a_{2}^{2}=\frac{\left(r_{2}+s_{2}\right)\left(h_{2}(x)\right)^{3}}{\left[(18 \alpha-6) \psi_{3}-(16 \alpha-8) \psi_{2}^{2}\right]\left(h_{2}(x)\right)^{2}-8 h_{3}(x)(2 \alpha-1)^{2}{\psi_{2}}^{2}} \tag{2.14}
\end{equation*}
$$

and by taking $h_{2}(x)=b x$ and $h_{3}(x)=b p x^{2}+q a$ in (2.14), it further yields

$$
\begin{equation*}
\left|a_{2}\right| \leq \frac{|b x| \sqrt{|b x|}}{\sqrt{\left|(9 \alpha-3) \psi_{3} b^{2} x^{2}-\left[(8 \alpha-4) b^{2} x^{2}+4\left(p x^{2} b+q a\right)(2 \alpha-1)^{2}\right] \psi_{2}{ }^{2}\right|}} . \tag{2.15}
\end{equation*}
$$

By subtracting (2.10) from (2.8) we get

$$
-6(3 \alpha-1) \psi_{3}\left(a_{2}^{2}-a_{3}\right)=\left(r_{2}-s_{2}\right) h_{2}(x)+\left(r_{1}^{2}-s_{1}^{2}\right) h_{3}(x)
$$

In view of (2.11), we obtain

$$
\begin{equation*}
a_{3}=\frac{\left(r_{2}-s_{2}\right) h_{2}(x)}{(18 \alpha-6) \psi_{3}}+a_{2}^{2} \tag{2.16}
\end{equation*}
$$

Then in view of (2.12), (2.16) becomes

$$
a_{3}=\frac{\left(r_{2}-s_{2}\right) h_{2}(x)}{(18 \alpha-6) \psi_{3}}+\frac{\left(h_{2}(x)\right)^{2}\left(r_{1}^{2}+s_{1}^{2}\right)}{8(2 \alpha-1)^{2} \psi_{2}^{2}}
$$

Applying (1.16), we deduce that

$$
\left|a_{3}\right| \leq \frac{|b x|}{(9 \alpha-3) \psi_{3}}+\frac{b^{2} x^{2}}{4(2 \alpha-1)^{2} \psi_{2}^{2}}
$$

From (2.16), for $\mu \in \mathbb{R}$, we write

$$
\begin{equation*}
a_{3}-\mu a_{2}^{2}=\frac{h_{2}(x)\left(r_{2}-s_{2}\right)}{(18 \alpha-6) \psi_{3}}+(1-\mu) a_{2}^{2} \tag{2.17}
\end{equation*}
$$

By substituting (2.14) in (2.17), we have

$$
\begin{align*}
a_{3}-\mu a_{2}^{2}= & \frac{h_{2}(x)\left(r_{2}-s_{2}\right)}{(18 \alpha-6) \psi_{3}} \\
& +\left(\frac{(1-\mu)\left(r_{2}+s_{2}\right)\left(h_{2}(x)\right)^{3}}{\left[(18 \alpha-6) \psi_{3}-(16 \alpha-8) \psi_{2}^{2}\right]\left(h_{2}(x)\right)^{2}-8 h_{3}(x)(2 \alpha-1)^{2} \psi_{2}^{2}}\right) \\
= & h_{2}(x)\left\{\left(\Lambda(\mu, x)+\frac{1}{(18 \alpha-6) \psi_{3}}\right) r_{2}\right. \\
& \left.+\left(\Lambda(\mu, x)-\frac{1}{(18 \alpha-6) \psi_{3}}\right) s_{2}\right\} \tag{2.18}
\end{align*}
$$

where

$$
\Lambda(\mu, x)=\frac{(1-\mu)\left[h_{2}(x)\right]^{2}}{\left[(18 \alpha-6) \psi_{3}-(16 \alpha-8) \psi_{2}^{2}\right]\left(h_{2}(x)\right)^{2}-8 h_{3}(x)(2 \alpha-1)^{2} \psi_{2}^{2}}
$$

Hence, we conclude that

$$
\left|a_{3}-\mu a_{2}^{2}\right| \leq \begin{cases}\frac{\left|h_{2}(x)\right|}{(9 \alpha-3) \psi_{3}} & ; 0 \leq|\Lambda(\mu, x)| \leq \frac{1}{(18 \alpha-6) \psi_{3}} \\ 2\left|h_{2}(x)\right||\Lambda(\mu, x)| & ;|\Lambda(\mu, x)| \geq \frac{1}{(18 \alpha-6) \psi_{3}}\end{cases}
$$

and in view of (1.16), it evidently completes the proof of Theorem 2.1.
Definition 2.2. A function $f \in \mathcal{A}$ of the form (1.1) belongs to the class $\mathcal{L}_{\sigma}(\lambda, \nu, q, x)$ for $\nu>0, \lambda>-1,0<q<1$, and $z, w \in \Delta$, if the following conditions are satisfied:

$$
\frac{1+\frac{z\left(\mathcal{N}_{\nu, q}^{\lambda} f(z)\right)^{\prime \prime}}{\left(\mathcal{N}_{\nu, q}^{\lambda} f(z)\right)^{\prime}}}{\frac{z\left(\mathcal{N}_{\nu, q}^{\lambda} f(z)\right)^{\prime}}{\mathcal{N}_{\nu, q}^{\lambda} f(z)}} \prec \Pi(x, z)+1-a
$$

and for $g$ the analytic extension (continuation) of $f^{-1}$ given by (1.2)

$$
\frac{1+\frac{w\left(\mathcal{N}_{\nu, q}^{\lambda} g(w)\right)^{\prime \prime}}{\left(\mathcal{N}_{\nu, q}^{\lambda} g(w)\right)^{\prime}}}{\frac{w\left(\mathcal{N}_{\nu, q}^{\lambda} g(w)\right)^{\prime}}{\mathcal{N}_{\nu, q}^{\lambda} g(w)}} \prec \Pi(x, w)+1-a,
$$

where the real constant $a$ is as in (1.16).
Remark 2.2. Putting $q \rightarrow 1^{-}$, we obtain that

$$
\lim _{q \rightarrow 1^{-}} \mathcal{L}_{\sigma}^{*}(\lambda, \nu, q, x)=: \mathcal{M}_{\sigma}(\lambda, \nu, x)
$$

where $f \in \sigma$,

$$
\frac{1+\frac{z\left(\mathcal{J}_{\nu}^{\lambda} f(z)\right)^{\prime \prime}}{\left(\mathcal{J}_{\nu}^{\lambda} f(z)\right)^{\prime}}}{\frac{z\left(\mathcal{J}_{\nu}^{\lambda} f(z)\right)^{\prime}}{\mathcal{J}_{\nu}^{\lambda} f(z)}} \prec \Pi(x, z)+1-a
$$

and for $g$ the analytic extension (continuation) of $f^{-1}$ given by (1.2)

$$
\frac{1+\frac{w\left(\mathcal{J}_{\nu}^{\lambda} g(w)\right)^{\prime \prime}}{\left(\mathcal{J}_{\nu}^{\lambda} g(w)\right)^{\prime}}}{\frac{w\left(\mathcal{J}_{\nu}^{\lambda} g(w)\right)^{\prime}}{\mathcal{J}_{\nu}^{\lambda} g(w)}} \prec \Pi(x, w)+1-a
$$

where the real constant $a$ is as in (1.16).
For functions in the class $\mathcal{L}_{\sigma}(\lambda, \nu, q, x)$, the following coefficient estimates and Fekete-Szegö inequality are obtained.
Theorem 2.2. Let $f(z)=z+\sum_{n=2}^{\infty} a_{n} z^{n}$ be in the class $\mathcal{L}_{\sigma}(\lambda, \nu, q, x)$. Then

$$
\left|a_{2}\right| \leq \frac{|b x| \sqrt{|b x|}}{\sqrt{\left|4 b^{2} x^{2} \psi_{3}-\left(4 b^{2} x^{2}+p x^{2} b+q a\right) \psi_{2}^{2}\right|}}, \quad \text { and } \quad\left|a_{3}\right| \leq \frac{|b x|}{4 \psi_{3}}+\frac{b^{2} x^{2}}{\psi_{2}^{2}}
$$

and for $\mu \in \mathbb{R}$

$$
\left|a_{3}-\mu a_{2}^{2}\right| \leq\left\{\begin{array}{c}
\frac{|b x|}{4 \psi_{3}} \\
\text { if } \quad|\mu-1| \leq \frac{\left|4 b^{2} x^{2} \psi_{3}-\left(4 b^{2} x^{2}+p x^{2} b+q a\right) \psi_{2}^{2}\right|}{4 b^{2} x^{2} \psi_{3}} \\
\frac{|b x|^{3}|\mu-1|}{\left|4 b^{2} x^{2} \psi_{3}-\left(4 b^{2} x^{2}+p x^{2} b+q a\right) \psi_{2}^{2}\right|} \\
\text { if } \quad|\mu-1| \geq \frac{\left|4 b^{2} x^{2} \psi_{3}-\left(4 b^{2} x^{2}+p x^{2} b+q a\right) \psi_{2}^{2}\right|}{4 b^{2} x^{2} \psi_{3}}
\end{array}\right.
$$

Proof. Let $f \in \mathcal{L}_{\sigma}(\lambda, \nu, q, x)$ be given by the Taylor-Maclaurin expansion (1.1). Then, there are analytic functions $r(z)$ and $s(w)$ such that

$$
r(0)=0 ; \quad s(0)=0, \quad\left|r_{n}\right|<1 \quad \text { and } \quad\left|s_{n}\right|<1 \quad(\forall z, w \in \Delta)
$$

and we can write

$$
\begin{equation*}
\frac{1+\frac{z\left(\mathcal{N}_{\nu, q}^{\lambda} f(z)\right)^{\prime \prime}}{\left(\mathcal{N}_{\nu, q}^{\lambda} f(z)\right)^{\prime}}}{\frac{z\left(\mathcal{N}_{\nu, q}^{\lambda} f(z)\right)^{\prime}}{\mathcal{N}_{\nu, q}^{\lambda} f(z)}}=\Pi(x, r(z))+1-a \tag{2.19}
\end{equation*}
$$

and

$$
\begin{equation*}
\frac{1+\frac{w\left(\mathcal{N}_{\nu, q}^{\lambda} g(w)\right)^{\prime \prime}}{\left(\mathcal{N}_{\nu, q}^{\lambda} g(w)\right)^{\prime}}}{\frac{w\left(\mathcal{N}_{\nu, q}^{\lambda} g(w)\right)^{\prime}}{\mathcal{N}_{\nu, q}^{\lambda} g(w)}}=\Pi(x, s(w))+1-a \tag{2.20}
\end{equation*}
$$

Equivalently,

$$
\begin{equation*}
\frac{1+\frac{z\left(\mathcal{N}_{\nu, q}^{\lambda} f(z)\right)^{\prime \prime}}{\left(\mathcal{N}_{\nu, q}^{\lambda} f(z)\right)^{\prime}}}{\frac{z\left(\mathcal{N}_{\nu, q}^{\lambda} f(z)\right)^{\prime}}{\mathcal{N}_{\nu, q}^{\lambda} f(z)}}=1+h_{1}(x)-a+h_{2}(x) r(z)+h_{3}(x)[r(z)]^{2}+\cdots \tag{2.21}
\end{equation*}
$$

and

$$
\begin{equation*}
\frac{1+\frac{w\left(\mathcal{N}_{\nu, q}^{\lambda} g(w)\right)^{\prime \prime}}{\left(\mathcal{N}_{\nu, q}^{\lambda} g(w)\right)^{\prime}}}{\frac{w\left(\mathcal{N}_{\nu, q}^{\lambda} g(w)\right)^{\prime}}{\mathcal{N}_{\nu, q}^{\lambda} g(w)}}=1+h_{1}(x)-a+h_{2}(x) s(w)+h_{3}(x)[s(w)]^{2}+\cdots . \tag{2.22}
\end{equation*}
$$

From (2.21) and (2.22) and in view of (1.17), we obtain

$$
\begin{equation*}
\frac{1+\frac{z\left(\mathcal{N}_{\nu, q}^{\lambda} f(z)\right)^{\prime \prime}}{\left(\mathcal{N}_{\nu, q}^{\lambda} f(z)\right)^{\prime}}}{\frac{z\left(\mathcal{N}_{\nu, q}^{\lambda} f(z)\right)^{\prime}}{\mathcal{N}_{\nu, q}^{\lambda} f(z)}}=1+h_{2}(x) r_{1} z+\left[h_{2}(x) r_{2}+h_{3}(x) r_{1}^{2}\right] z^{2}+\cdots \tag{2.23}
\end{equation*}
$$

and

$$
\begin{equation*}
\frac{1+\frac{w\left(\mathcal{N}_{\nu, q}^{\lambda} g(w)\right)^{\prime \prime}}{\left(\mathcal{N}_{\nu, q}^{\lambda} g(w)\right)^{\prime}}}{\frac{w\left(\mathcal{N}_{\nu, q}^{\lambda} g(w)\right)^{\prime}}{\mathcal{N}_{\nu, q}^{\lambda} g(w)}}=1+h_{2}(x) s_{1} w+\left[h_{2}(x) s_{2}+h_{3}(x) s_{1}^{2}\right] w^{2}+\cdots . \tag{2.24}
\end{equation*}
$$

If

$$
r(z)=\sum_{n=1}^{\infty} r_{n} z^{n} \quad \text { and } \quad s(w)=\sum_{n=1}^{\infty} s_{n} w^{n},
$$

then it is well known that

$$
\left|r_{n}\right| \leq 1 \quad \text { and } \quad\left|s_{n}\right| \leq 1 \quad(n \in \mathbb{N})
$$

Thus upon comparing the corresponding coefficients in (2.23) and (2.24), we have

$$
\begin{align*}
\psi_{2} a_{2} & =h_{2}(x) r_{1}  \tag{2.25}\\
4\left(a_{3} \psi_{3}-a_{2}^{2} \psi_{2}^{2}\right) & =h_{2}(x) r_{2}+h_{3}(x) r_{1}^{2}  \tag{2.26}\\
-\psi_{2} a_{2} & =h_{2}(x) s_{1} \tag{2.27}
\end{align*}
$$

and

$$
\begin{equation*}
\left(8 a_{2}^{2}-4 a_{3}\right) \psi_{3}-4 a_{2}^{2} \psi_{2}^{2}=h_{2}(x) s_{2}+h_{3}(x) s_{1}^{2} . \tag{2.28}
\end{equation*}
$$

The results of this theorem now follow from (2.25)-(2.28) by applying the procedure as in Theorem 2.1 with respect to (2.7)-(2.10).

Definition 2.3. A function $f \in \mathcal{A}$ of the form (1.1) belongs to the class $\mathcal{P}_{\sigma}(\gamma, \lambda$, $\nu, q, x)$ for $0 \leq \gamma \leq 1, \nu>0, \lambda>-1,0<q<1$, and $z, w \in \Delta$, if the following conditions are satisfied:

$$
\frac{z\left(\mathcal{N}_{\nu, q}^{\lambda} f(z)\right)^{\prime}+\gamma z^{2}\left(\mathcal{N}_{\nu, q}^{\lambda} f(z)\right)^{\prime \prime}}{(1-\gamma)\left(\mathcal{N}_{\nu, q}^{\lambda} f(z)\right)+\gamma z\left(\mathcal{N}_{\nu, q}^{\lambda} f(z)\right)^{\prime}} \prec \Pi(x, z)+1-a
$$

and for $g$ the analytic extension (continuation) of $f^{-1}$ given by (1.2)

$$
\frac{w\left(\mathcal{N}_{\nu, q}^{\lambda} g(w)\right)^{\prime}+\gamma w^{2}\left(\mathcal{N}_{\nu, q}^{\lambda} g(w)\right)^{\prime \prime}}{(1-\gamma)\left(\mathcal{N}_{\nu, q}^{\lambda} g(w)\right)+\gamma w\left(\mathcal{N}_{\nu, q}^{\lambda} g(w)\right)^{\prime}} \prec \Pi(x, w)+1-a
$$

where the real constant $a$ is as in (1.16).
Remark 2.3. Putting $q \rightarrow 1^{-}$, we obtain that

$$
\lim _{q \rightarrow 1^{-}} \mathcal{P}_{\sigma}(\gamma, \lambda, \nu, q, x)=: \mathcal{P}_{\sigma}(\gamma, \lambda, \nu, x)
$$

where $f \in \sigma$,

$$
\frac{z\left(\mathcal{J}_{\nu}^{\lambda} f(z)\right)^{\prime}+\gamma z^{2}\left(\mathcal{J}_{\nu}^{\lambda} f(z)\right)^{\prime \prime}}{(1-\gamma)\left(\mathcal{J}_{\nu}^{\lambda} f(z)\right)+\gamma z\left(\mathcal{J}_{\nu}^{\lambda} f(z)\right)^{\prime}} \prec \Pi(x, z)+1-a
$$

and for $g$ the analytic extension (continuation) of $f^{-1}$ given by (1.2)

$$
\frac{w\left(\mathcal{J}_{\nu}^{\lambda} g(w)\right)^{\prime}+\gamma w^{2}\left(\mathcal{J}_{\nu}^{\lambda} g(w)\right)^{\prime \prime}}{(1-\gamma)\left(\mathcal{J}_{\nu}^{\lambda} g(w)\right)+\gamma w\left(\mathcal{J}_{\nu}^{\lambda} g(w)\right)^{\prime}} \prec \Pi(x, w)+1-a,
$$

For functions in the class $\mathcal{P}_{\sigma}(\gamma, \lambda, \nu, q, x)$, the following coefficient estimates and Fekete-Szegö inequality are obtained.
Theorem 2.3. Let $f(z)=z+\sum_{n=2}^{\infty} a_{n} z^{n}$ be in the class $\mathcal{P}_{\sigma}(\gamma, \lambda, \nu, q, x)$. Then

$$
\begin{aligned}
& \left|a_{2}\right| \leq \frac{|b x| \sqrt{|b x|}}{\sqrt{\left|2 b^{2} x^{2}(2 \gamma+1) \psi_{3}-(1+\gamma)^{2}\left(b^{2} x^{2}+p x^{2} b+q a\right) \psi_{2}{ }^{2}\right|}}, \\
& \left|a_{3}\right| \leq \frac{|b x|}{2 \psi_{3}(2 \gamma+1)}+\frac{b^{2} x^{2}}{(1+\gamma)^{2} \psi_{2}{ }^{2}}
\end{aligned}
$$

and for $\mu \in \mathbb{R}$

$$
\left|a_{3}-\mu a_{2}^{2}\right| \leq\left\{\begin{array}{l}
\frac{|b x|}{2 \psi_{3}(2 \gamma+1)} \\
\text { if } \quad|\mu-1| \leq \frac{\left|2 b^{2} x^{2}(2 \gamma+1) \psi_{3}-(1+\gamma)^{2}\left(b^{2} x^{2}+p x^{2} b+q a\right) \psi_{2}^{2}\right|}{2 b^{2} x^{2} \psi_{3}(2 \gamma+1)} \\
\frac{|b x|^{3}|\mu-1|}{\left|2 b^{2} x^{2}(2 \gamma+1) \psi_{3}-(1+\gamma)^{2}\left(b^{2} x^{2}+p x^{2} b+q a\right) \psi_{2}^{2}\right|} \\
\text { if } \quad|\mu-1| \geq \frac{\left|2 b^{2} x^{2}(2 \gamma+1) \psi_{3}-(1+\gamma)^{2}\left(b^{2} x^{2}+p x^{2} b+q a\right) \psi_{2}^{2}\right|}{2 b^{2} x^{2} \psi_{3}(2 \gamma+1)} .
\end{array}\right.
$$

Proof. Let $f \in \mathcal{P}_{\sigma}(\gamma, \lambda, \nu, q, x)$ be given by the Taylor-Maclaurin expansion (1.1). Then, there are analytic functions $r(z)$ and $s(w)$ such that

$$
r(0)=0 ; \quad s(0)=0, \quad\left|r_{n}\right|<1 \quad \text { and } \quad\left|s_{n}\right|<1 \quad(\forall z, w \in \Delta),
$$

and we can write

$$
\begin{equation*}
\frac{z\left(\mathcal{N}_{\nu, q}^{\lambda} f(z)\right)^{\prime}+\gamma z^{2}\left(\mathcal{N}_{\nu, q}^{\lambda} f(z)\right)^{\prime \prime}}{(1-\gamma)\left(\mathcal{N}_{\nu, q}^{\lambda} f(z)\right)+\gamma z\left(\mathcal{N}_{\nu, q}^{\lambda} f(z)\right)^{\prime}}=\Pi(x, r(z))+1-a \tag{2.29}
\end{equation*}
$$

and

$$
\begin{equation*}
\frac{w\left(\mathcal{N}_{\nu, q}^{\lambda} g(w)\right)^{\prime}+\gamma w^{2}\left(\mathcal{N}_{\nu, q}^{\lambda} g(w)\right)^{\prime \prime}}{(1-\gamma)\left(\mathcal{N}_{\nu, q}^{\lambda} g(w)\right)+\gamma w\left(\mathcal{N}_{\nu, q}^{\lambda} g(w)\right)^{\prime}}=\Pi(x, s(w))+1-a \tag{2.30}
\end{equation*}
$$

Equivalently,

$$
\begin{align*}
& \frac{z\left(\mathcal{N}_{\nu, q}^{\lambda} f(z)\right)^{\prime}+\gamma z^{2}\left(\mathcal{N}_{\nu, q}^{\lambda} f(z)\right)^{\prime \prime}}{(1-\gamma)\left(\mathcal{N}_{\nu, q}^{\lambda} f(z)\right)+\gamma z\left(\mathcal{N}_{\nu, q}^{\lambda} f(z)\right)^{\prime}} \\
& \quad=1+h_{1}(x)-a+h_{2}(x) r(z)+h_{3}(x)[r(z)]^{2}+\cdots \tag{2.31}
\end{align*}
$$

and

$$
\begin{align*}
& \frac{w\left(\mathcal{N}_{\nu, q}^{\lambda} g(w)\right)^{\prime}+\gamma w^{2}\left(\mathcal{N}_{\nu, q}^{\lambda} g(w)\right)^{\prime \prime}}{(1-\gamma)\left(\mathcal{N}_{\nu, q}^{\lambda} g(w)\right)+\gamma w\left(\mathcal{N}_{\nu, q}^{\lambda} g(w)\right)^{\prime}} \\
& \quad=1+h_{1}(x)-a+h_{2}(x) s(w)+h_{3}(x)[s(w)]^{2}+\cdots \tag{2.32}
\end{align*}
$$

From (2.31) and (2.32) and in view of (1.17), we obtain

$$
\begin{align*}
& \frac{z\left(\mathcal{N}_{\nu, q}^{\lambda} f(z)\right)^{\prime}+\gamma z^{2}\left(\mathcal{N}_{\nu, q}^{\lambda} f(z)\right)^{\prime \prime}}{(1-\gamma)\left(\mathcal{N}_{\nu, q}^{\lambda} f(z)\right)+\gamma z\left(\mathcal{N}_{\nu, q}^{\lambda} f(z)\right)^{\prime}} \\
& \quad=1+h_{2}(x) r_{1} z+\left[h_{2}(x) r_{2}+h_{3}(x) r_{1}^{2}\right] z^{2}+\cdots \tag{2.33}
\end{align*}
$$

and

$$
\begin{align*}
& \frac{w\left(\mathcal{N}_{\nu, q}^{\lambda} g(w)\right)^{\prime}+\gamma w^{2}\left(\mathcal{N}_{\nu, q}^{\lambda} g(w)\right)^{\prime \prime}}{(1-\gamma)\left(\mathcal{N}_{\nu, q}^{\lambda} g(w)\right)+\gamma w\left(\mathcal{N}_{\nu, q}^{\lambda} g(w)\right)^{\prime}} \\
& \quad=1+h_{2}(x) s_{1} w+\left[h_{2}(x) s_{2}+h_{3}(x) s_{1}^{2}\right] w^{2}+\cdots \tag{2.34}
\end{align*}
$$

If

$$
r(z)=\sum_{n=1}^{\infty} r_{n} z^{n} \quad \text { and } \quad s(w)=\sum_{n=1}^{\infty} s_{n} w^{n}
$$

then it is well known that

$$
\left|r_{n}\right| \leq 1 \quad \text { and } \quad\left|s_{n}\right| \leq 1 \quad(n \in \mathbb{N})
$$

Thus upon comparing the corresponding coefficients in (2.33) and (2.34), we have

$$
\begin{equation*}
(1+\gamma) \psi_{2} a_{2}=h_{2}(x) r_{1} \tag{2.35}
\end{equation*}
$$

$$
\begin{gather*}
2(1+2 \gamma) \psi_{3} a_{3}-(1+\gamma)^{2} \psi_{2}^{2} a_{2}^{2}=h_{2}(x) r_{2}+h_{3}(x) r_{1}^{2}  \tag{2.36}\\
-(1+\gamma) \psi_{2} a_{2}=h_{2}(x) s_{1} \tag{2.37}
\end{gather*}
$$

and

$$
\begin{equation*}
\left((8 \gamma+4) \psi_{3}-\psi_{2}^{2}(1+\gamma)^{2}\right) a_{2}^{2}-2 a_{3}(2 \gamma+1) \psi_{3}=h_{3}(x) s_{1}^{2}+h_{2}(x) s_{2} \tag{2.38}
\end{equation*}
$$

The results of this theorem now follow from (2.35)-(2.38) by applying the procedure as in Theorem 2.1 with respect to (2.7)-(2.10).

## 3. Corollaries and Consequences

Taking $\gamma=0$ in Theorem (2.3), we have following corollary.
Corollary 3.1. Let $f(z)=z+\sum_{n=2}^{\infty} a_{n} z^{n}$ be in the class $\mathcal{S}_{\sigma}(\lambda, \nu, q, x)$. Then

$$
\left|a_{2}\right| \leq \frac{|b x| \sqrt{|b x|}}{\sqrt{\left|2 b^{2} x^{2} \psi_{3}-\left(b^{2} x^{2}+p x^{2} b+q a\right) \psi_{2}{ }^{2}\right|}}, \quad\left|a_{3}\right| \leq \frac{|b x|}{2 \psi_{3}}+\frac{b^{2} x^{2}}{\psi_{2}{ }^{2}}
$$

and for $\mu \in \mathbb{R}$

$$
\left|a_{3}-\mu a_{2}^{2}\right| \leq\left\{\begin{array}{c}
\frac{|b x|}{2 \psi_{3}} \\
\text { if } \quad|\mu-1| \leq \frac{\left|2 b^{2} x^{2} \psi_{3}-\left(b^{2} x^{2}+p x^{2} b+q a\right) \psi_{2}^{2}\right|}{2 b^{2} x^{2} \psi_{3}} \\
\frac{|b x|^{3}|\mu-1|}{\left|2 b^{2} x^{2} \psi_{3}-\left(b^{2} x^{2}+p x^{2} b+q a\right) \psi_{2}^{2}\right|} \\
\text { if } \quad|\mu-1| \geq \frac{\left|2 b^{2} x^{2} \psi_{3}-\left(b^{2} x^{2}+p x^{2} b+q a\right) \psi_{2}^{2}\right|}{2 b^{2} x^{2} \psi_{3}} .
\end{array}\right.
$$

Taking $\alpha=1$ in Theorem 2.1 or $\gamma=1$ in Theorem 2.3, we have following corollary.
Corollary 3.2. Let $f(z)=z+\sum_{n=2}^{\infty} a_{n} z^{n}$ be in the class $\mathcal{K}_{\sigma}(\lambda, \nu, q, x)$. Then

$$
\left|a_{2}\right| \leq \frac{|b x| \sqrt{|b x|}}{\sqrt{\left|6 b^{2} x^{2} \psi_{3}-4\left(b^{2} x^{2}+p x^{2} b+q a\right){\psi_{2}^{2}}^{2}\right|}}, \quad\left|a_{3}\right| \leq \frac{|b x|}{6 \psi_{3}}+\frac{b^{2} x^{2}}{4 \psi_{2}{ }^{2}}
$$

and for $\mu \in \mathbb{R}$

$$
\left|a_{3}-\mu a_{2}^{2}\right| \leq\left\{\begin{array}{c}
\frac{|b x|}{6 \psi_{3}} \\
\text { if } \quad|\mu-1| \leq \frac{\left|3 b^{2} x^{2} \psi_{3}-2\left(b^{2} x^{2}+p x^{2} b+q a\right) \psi_{2}^{2}\right|}{3 b^{2} x^{2} \psi_{3}} \\
\frac{|b x|^{3}|\mu-1|}{\left|6 b^{2} x^{2} \psi_{3}-4\left(b^{2} x^{2}+p x^{2} b+q a\right) \psi_{2}^{2}\right|} \\
\text { if } \quad|\mu-1| \geq \frac{\left|3 b^{2} x^{2} \psi_{3}-2\left(b^{2} x^{2}+p x^{2} b+q a\right) \psi_{2}^{2}\right|}{3 b^{2} x^{2} \psi_{3}}
\end{array} .\right.
$$

## 4. Conclusion

One could find initial coefficient estimates for the classes defined in Remarks 2.1, 2.2 and 2.3. We leave those to interested readers.

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